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INVESTIGATION ON ELECTRONIC AND THERMOELECTRIC PROPERTIES OF (P, As, Sb) DOPED ZrCoBi

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Since the last decade, the half-Heusler (HH) compounds have taken an important place in the field of the condensed matter physics research. The multiplicity of substitutions of transition elements at the crystallographic sites X, Y and (III-V) elements at the Z sites, gives to the HH alloys a multitudes of remarkable properties. In the present study, we examined the structural, electronic and thermoelectric properties of ZrCoBi_{0.75}Z_{0.25} (Z = P, As, Sb) using density functional theory (DFT). The computations have been done parallel to the full potential linearized augmented plane wave (FP-LAPW) method as implemented in the WIEN2k code. The thermoelectrically properties were predicted via the semi-classical Boltzmann transport theory, as performed in Boltztrap code. The obtained results for the band structure and densities of states confirm the semiconductor (SC) nature of the three compounds with an indirect band gap, which is around 1eV. The main thermoelectric parameters such as Seebeck coefficient, thermal conductivity, electrical conductivity and figure of merit were estimated for temperatures ranging from zero to 1200K. The positive values of Seebeck coefficient (S) confirm that the $ZrCoBi_{0.75}Z_{0.25}$ (x = 0 and 0.25) are a p-type SC. At the ambient temperature, $ZrCoBi_{0.75}P_{0.25}$ exhibit the large S value of 289 µV/K, which constitutes an improvement of 22% than the undoped ZrCoBi, and show also a reduction of 54% in thermal conductivity (κ/τ). The undoped ZrCoBi has the lowest ZT value at all temperatures and by substituting bismuth atom by one of the sp elements (P, As, Sb), a simultaneous improvement in κ/τ and S have led to maximum figure of merit (ZT) values of about 0.84 obtained at 1200 K for the three-doped compounds.

KEYWORDS: First-principle, GGA, Doped-semiconductor, Seebeck coefficient, figure of merit

The half-Heusler (HH) semiconductors have attracted substantial interest of researchers, they are considered as alternative and green energy sources because they have the capacity to transform waste heat into electric energy. A significant change in their thermoelectric properties can be observed by the incorporation of a dopant at various concentration rates. Many researches on the thermoelectric properties of HH compounds show moderately high ZT of ≈1.0 in the mid to high temperature ranges (700–1000 K) [1] which makes them as potential candidates for high efficiency power generation. Recently, several experimental [2-4] and theoretical [5-7] studies have been focused on the doped HH alloys, these last can be utilized in a broad temperature region and a good part of their raw materials are nontoxic, abundant and cheap [8-11]. The single doping remains one of the effective process, which adjust the electrical power factor and reduce the thermal conductivity [12]. C.C. Hsu et al., [13] have studied the ZrFe_xCo_{1.x}Sb HH and they show that for x=0.2, the maximum ZT value of 0.036 was obtained at 900K, they attribute this to reduction of thermal conductivity because the point defects, mass alterations and strain field effects induced by doping vigorously scatter the thermal phonons. Y.Lei et al. [14] have explored the impact of Sb doping on thermoelectric properties of TiNiSn HH, they noted that the ZT value is moderately improved from 0.30 for undoped alloy to 0.44 for a dopant concentration of 3%. R. Akram et al [15] find that the incorporation of 1.5% of Sb in Hf_{0.25}Zr_{0.75}NiSn HH enhanced the electrical properties by increasing the carrier concentration and reducing the thermal conductivity. A maximal ZT value of 0.83 was achieved at 923 K, which constitute 67% of improvement compared to the undoped sample value. In another work [16], the La_xZr_{1-x}NiSn compounds, was synthetized by induction melting combined with plasma-activated sintering, the x-ray diffraction confirms that for all x-concentration of La, the structures remain cubic with space group F-43m. A particular drop in thermal conductivity was found following doping by La, thus leading to ZT of 0.53 at 923 K, which represent an enhancement of about 37% compared to the pristine sample. In the present work, we studied the electronic and thermoelectric properties of sp elements (P, As, Sb)-doped ZrCoBi half-Heusler alloy. A brief commentary on the band structure and densities of states are given. A particular emphasis was placed on the thermoelectric properties investigated by the semi-classical Boltzmann theory.

COMPUTATIONAL METHOD

The computations are performed by the full potential linearized augmented plane wave (FP-LAPW) method [17-18] based on DFT [19] as implemented in Wien2K package [20]. The generalized gradient approximation (GGA) [21] is employed for the exchange-correlation potential. The muffin-tin sphere radii RMT were chosen as equal to 1.85, 2.1 and 2.5 a.u for the Zr, Co and (Bi / P, As, Sb) atoms respectively. For convergence, the energy cut off = -6.0 Ry, Gmax =12, lmax =10 and RMT*K_{max} = 7 are used. The 22×22×22 k-point mesh is used for electronic properties calculation. The self-consistency is achieved up to 0.0001 Ry. The space group of cubic ZrCoBi is 216 (F-43m). According to reference [22], the compound is stable in paramagnetic phase and the atomic coordinates are (0, 0, 0), (0.25, 0.25, 0.25) and (0.5, 0.5, 0.5) for Zr, Co and Bi respectively. With the aim of creating 25 % concentration of doping in ZrCoBi lattice, the supercell $1\times1\times1$ containing twelve atoms is constructed. The dopant elements (Phosphor, Arsenic and Antimony) was substituted at the Bismuth site. The space group number of this generated super cell is 215 (P-43m) which contains 4 Zr atoms, 4 Co atoms and 4 Bi atoms. For the four Bi atom, the one who is situated at the (0, 0, 0) position was replaced with (P/As/Sb) atom without disturbing the other eleven atoms (4 Zr, 4 Co and 3 Bi) in the supercell. The semi-classical Boltzmann approach [23] as given in the BoltzTraP code associated to a fine grid mesh of $(46\times46\times46)$ were employed to investigate the thermoelectric response of compounds doped and undoped.

RESULTS AND DISCUSSION

Structural and electronic properties

The crystal structure of the perfect ZrCoBi HH and the doped $ZrCoBi_{0.75}Z_{0.25}$ (Z = P, As, Sb) are plotted by the CrystalMaker 2.7 software [24]. The undoped structure of ZrCoBi (Fig.1a) is formed by three interpenetrating fcc sublattices, which are occupied by Zr, Co and Bi elements. Its lattice constant obtained from optimization is in good agreement with the value of 6.22 Å, achieved by G. Surucu [22] using the GGA-PBE.

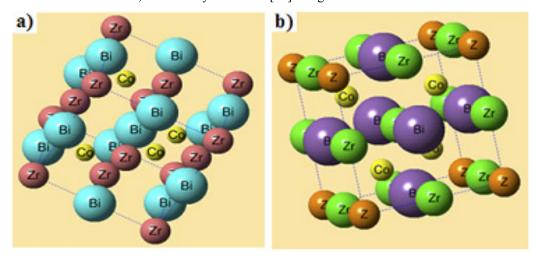


Figure 1. a) - Crystal structure of half-Heusler semiconductor ZrCoBi, b) - 12 atom simple cubic model of ZrCoBi_{0.75}Z_{0.25} (Z = P, As, Sb) structure.

The crystal structures of $ZrCoBi_{0.75}Z_{0.25}$ (Z = P, As, Sb) are cubic with the space group of P-43m; the structure can be regarded as four interpenetrating fcc sublattices (Fig. 1b). The calculated structural properties such as lattice constants a, bulk modulus B, first derivatives B' and equilibrium energies Emin are recapitulated in Table 1.

Table 1. Calculated lattice constant $a_0(A^\circ)$, bulk modulus B(GPa), derivative of the bulk modulus, B'(GPa), minimum total energy per unit cell E_{min} (Ry) and Gap energy (eV)

Compounds	Space group	A (Å)	B (GPa)	В'	Emin (Ry)	Gap (eV)
Undoped - ZrCoBi	216 F-43m	6.23	120.31	4.94	-53148.552763	0.990
ZrCoBi _{0.75} P _{0.25}		6.12	131.34	5.23	-170115.352819	1.018
ZrCoBi _{0.75} As _{0.25}	215 P-43m	6.14	128.07	5.10	-173953.361417	1.008
ZrCoBio 75Sbo 25		6.19	124.95	4.95	-182398.455731	0.996

The lattice constants for ZrCoBi_{0.75}P_{0.25}, ZrCoBi_{0.75}As_{0.25}, ZrCoBi_{0.75}As_{0.25}, ZrCoBi_{0.75}Sb_{0.25}, show increase from P to Sb dopant. The three doping elements are electronegative and their atomic radii in increasing order is P, As and Sb. Fig. 2(a–d) show the calculated electronic band structure with high symmetric Brillouin zone W-L-Γ-X-W-K for undoped ZrCoBi and ZrCoBi_{0.75}Z_{0.25} (Z = P, As, Sb) Half-Heusler using the GGA approach. For the pristine case, the Fermi level lies inside the forbidden gap and the valance band maxima (VBM) and conduction band minima (CBM) occur at the L and X points, thus generating an indirect band gap of about 0.99 eV, which confirms the semiconducting nature of compound. The introduction of sp elements (P, As, Sb) ions into ZrCoBi shifts the bottom point of the conduction band towards a higher energy than the undoped ZrCoBi, thus increasing the band gap. The values to 0.996 eV, 1.008 eV and 1.018 eV are obtained for ZrCoBi_{0.75}Sb_{0.25}, ZrCoBi_{0.75}As_{0.25} and ZrCoBi_{0.75}P_{0.25} respectively. The VBM and the CBM are located at the L and Γ points, which indicates the change in momentum of the electrons during their transitions from VBM to CBM. Consequently, all the ZrCoBi_{0.75}Z_{0.25} (Z = P, As, Sb) compounds are indirect band gap semiconductors. In order to study the positioning of the orbital's in the electronic band structure as well as the electrons involving in the shaping of the band gap, the total anacd partial density of states (TDOS/PDOSs) of pure and sp element (P, As, and Sb) doped ZrCoBi are investigated between -4 and 4 eV Fig. 3(a-d). The dashed line displays the Fermi energy level (EF). We can see that the

top of the VB is mainly determined by the hybridization of d-Zr, d-Co, P-Bi and p-P states, whereas the bottom of the CB is formed primarily by the d-states of Co and Zr. From [- 4 to - 2eV], the anion d-Zr, d-Co, P-Bi and p-P states depict the lowest bands, these states are separated the main valence state (d-Co) by a small gap. The (PDOS) in this energy range is more pronounced for P-doped ZrCoBi than for As-doped ZrCoBi or Sb-doped ZrCoBi. The density of states near the Fermi level for the ZrCoBi_{0.75}Z_{0.25} (Z = P, As, Sb) HH, can be attributed to d-d orbitals hybridization between transition metals (Zr/Co). The p-orbitals of the Z atom contribute only slightly to the electronic properties. The offset of the valence bands toward the low energies is more larger for ZrCoBi_{0.75}Z_{0.25} (Z = P, As, Sb) than that for undoped-ZrCoBi, while the conduction bands remain almost in the same energy range (Fig. 3a). We observe that with increasing in atomic number of dopant element a shift of the main peak (d-Co) towards higher energies occurs [P (-1.22 eV)/As (-1.08 eV)/Sb (-1.04 eV)].

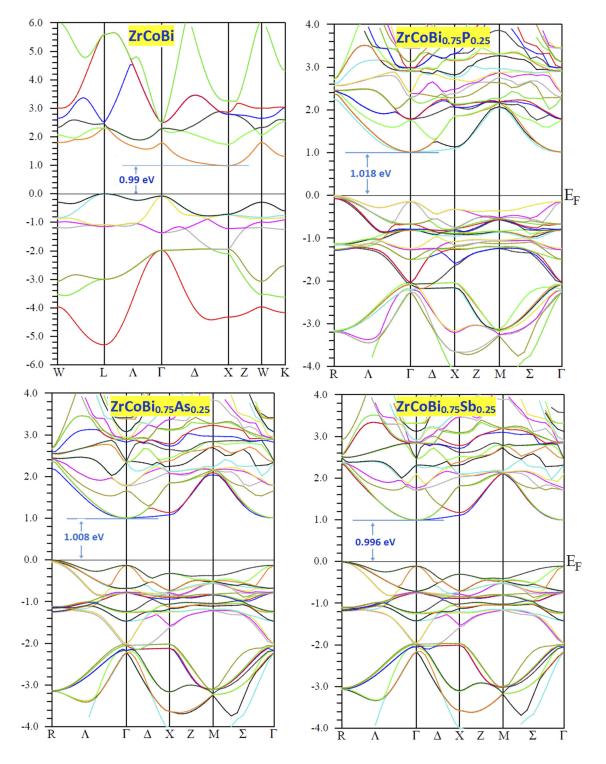


Figure 2. Band structure for a) undoped ZrCoBi, b) ZrCoBio.25Po.25, c) ZrCoBio.25Aso.25 and d) ZrCoBio.25Sbo.25.

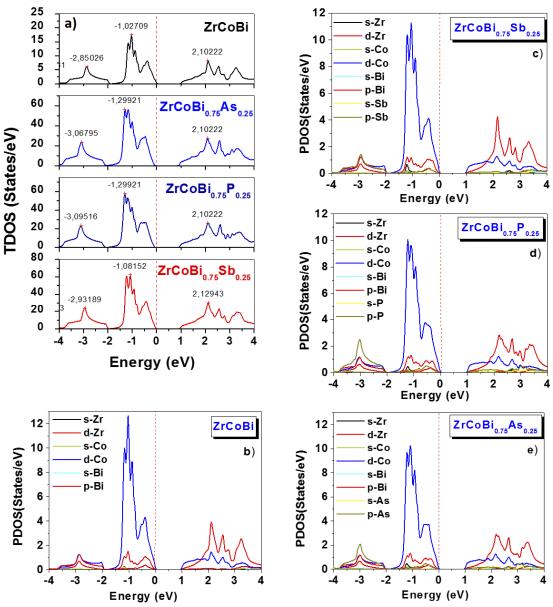


Figure 3. a) - Total density of states (TDOS), b-d) - Partial densities of states (PDOS)

Thermoelectric properties

The appropriate choice of the doping as well as its concentration remain among the principal factors which maximize the ZT of semiconductors and which make them suitable for the thermoelectric applications. The transport properties will be computed with BoltzTraP code under a constant relaxation time approximation of the charge carriers. The main thermoelectric parameters, such as Seebeck coefficient (S), electrical conductivity (σ/τ), thermal conductivity (κ/τ) and figure of merit (ZT) will be investigated as function of temperature [25-26]. The temperature reliance of S is shown in Fig. 4c. It is observed that beyond temperature of 400 K, the Seebeck coefficient (S) decrease steadily with the increase of temperature for the $ZrCoBi_{0.75}Z_{0.25}(Z = P, As, Sb)$ HH. This limitation of S at high temperature can be attributed to the bipolar effects which occur for wide-bandgap semiconductors and which result in an increase of thermal conductivity [27]. Among the three doped compounds, the P-doped ZrCoBi exhibit the large Seebeck coefficient of 289 μV/K between 300 to 400 K, this obtained value of S has been improved by 22% than the undoped ZrCoBi. With the increased of the atomic mass of the doping element, a diminution of S was observed. The positive values of Seebeck coefficient confirm that the $ZrCoBi_{0.75}Z_{0.25}$ (Z = P, As, Sb) are p-type semiconductors. The obtained values of S around room temperature are larger to those reported by several doped and undoped HH compounds such TiCoSn_xSb_{1-x} [28], FeV_{1-x}Hf_xSb [29], PdZrGe [30] ScRhTe [31]. Our compounds are good thermoelectric materials (TM), because according to J.W. Sharp [32], the best TM are highly doped semiconductors with a Seebeck coefficient located in the domain $150 - 250 \,\mu\text{V K}^{-1}$. The ability of the ZrCoBi_{0.75}Z_{0.25} (Z = P, As, Sb) HH to transfer heat is investigated in this section by the computation of its thermal conductivity (κ/τ). Fig. 4b shows temperature dependence of κ/τ of doped and undoped ZrCoBi.

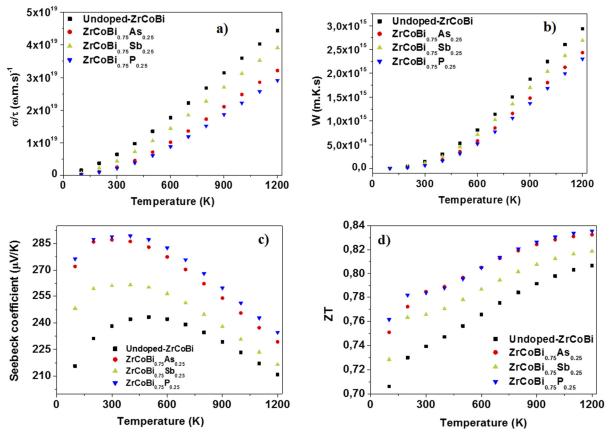


Figure 4. Thermoelectric properties versus temperature a) Electrical conductivity b) Thermal conductivity c) Seebeck coefficient and d) figure of merit

We can see that doping reliance of the κ/τ showed a decreasing value comparatively to the undoped ZrCoBi, which is advantageous for improving the ZT value. After 400 K, the κ/τ , which is directly allied with electrical conductivity, augments with temperature in a very linear trend and the slope of the κ/τ curves increases with increasing the Z-atomic number. The κ/τ which is also influenced by the evolution of charge-carrier concentration, show at 300 K a value of about 5.90.10¹³ W/m.K.s which is comparable with that of Pb₇Bi₄Se₁₃ [33], CdSe, CdTe [34]. In comparison with undoped ZrCoBi, a reduction of nearly 23% in the κ/τ was obtained at 1200 K by the substitution of the Bismuth site by the Phosphor. Around the ambient temperature, the κ/τ is weak then it sudden increases with the increasing in temperature this is due to electron-phonon scattering [35]. With an increase from 5.91.10¹² mKs at 100 K to 2.69.10¹⁵ mKs at 1200 K, and with value of 6.50.1013 W/mKs at room temperature (300 K), the ZrCoBi_{0.75}Sb_{0.25} constitutes a good thermal conductor. The effectiveness of thermoelectric devices requires a sufficiently high electrical conductivity (σ/τ) to achieve a high figure of merit (ZT). Fig. 4(a) shows the temperature dependence of σ/τ for the ZrCoBi_{0.75}Z_{0.25} compounds. The σ/τ displays exponential growth with increasing temperature due to the increase in the number of carriers. The ZrCoBi_{0.75}Sb_{0.25} exhibits the highest value of σ/τ comparatively to ZrCoBi_{0.75}As_{0.25} and ZrCoBi_{0.75}P_{0.25}, because the antimony shows the greater atomic weight to that of arsenic and phosphour. The increase in atomic weight augments the number of electronic layer and nucleus grip on the valence electrons then becomes weak. Several electrons are pushed toward the conduction band leading to improve the σ/τ values [23]. At room temperature, the σ/τ is of 4.41.10¹⁸, 2.52.10¹⁸ and $2.09.10^{18}$ (Ω .m.s)⁻¹ for ZrCoBi_{0.75}Sb_{0.25}, ZrCoBi_{0.75}As_{0.25} and ZrCoBi_{0.75}P_{0.25} respectively. These values increase with the temperature and reaches the maximum of $3.91.10^{19}$, $3.21.10^{19}$ and $2.91.10^{19}$ (Ω .m.s)⁻¹ at 1200 K. The increasing trend of σ/τ displayed for doped ZrCoBi and undoped ZrCoBi is in agreement with previous studies done on the semiconductors compounds [36-37]. The calculated values of figure of merit versus temperature are presented in Fig. 4d. There is a continuous increase in ZT with increasing temperature. The undoped ZrCoBi has the lowest ZT value at all temperatures and by substituting bismuth atom by one of the sp elements (P, As, Sb), an increase of the ZT was observed. A maximum ZT = 0.85 was obtained at 1200 K with a 25% P-substitution at the Bi-site. At low temperature (100 K), the deviation of ZT between the doped and undoped structures is much more marked. The high values of the Seebeck coefficient as well as the electrical conductivity of ZrCoBi_{1-x}Z_x alloys did not translate into high values of ZT because of rather high thermal conductivity. The difference in mass and radius between Bi and Z are responsible for the high thermal conductivity. However, others adjustment remains necessary for more improvement let's quote the concentration of the dopant element which should be optimized to decrease more the density of transport electrons also the use of second substituting atom (co-doping) with larger mass can reduce the thermal conductivity [38]. Despite that, the figure of merit and Seebeck coefficients values obtained and which have been widely improved by doping we can conclude that ours doped compounds have the capacity to transform waste heat into electric energy and may be also used in various thermoelectrically devices.

CONCLUSION

By using the first-principle density-functional theory within the supercell approach, we have studied the $ZrCoBi_{0.75}Z_{0.25}$ (Z=P, As, Sb) half-Heusler, a particular emphasis was placed on the thermoelectric properties. The calculated band structures reveal that the doped and undoped-ZrCoBi are indirect band gap semiconductors. The substitution of the Bi site by P, As and Sb is found effective in decreasing thermal conductivity and increasing the Seebeck coefficient and figure of merit value. Comparatively to undoped-ZrCoBi, the $ZrCoBi_{0.75}P_{0.25}$ show improvements of 22% and 54% of its Seebeck coefficient and thermal conductivity respectively. An enhancement of ZT values of 7.8% and 6.1% have been observed at 100K and 300K between the undoped-ZrCoBi and P-doped ZrCoBi half-Heusler. In order to further boost the performance of these alloys, future work will focus on single doping with very low dopant concentration as well as on co-doping.

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ДОСЛІДЖЕННЯ ЕЛЕКТРОННИХ ТЕРМОЕЛЕКТРИЧНИХ ВЛАСТИВОСТЕЙ ZrCoBi ЛЕГОВАНОГО (P, As, Sb)

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Протягом останнього десятиріччя напів-Гайслерові сполуки (НН) займають важливе місце в галузі досліджень фізики конденсованої речовини. Численність заміщень перехідних елементів на кристалографічних ділянках X, Y та (III-V) елементів на Z-ділянках надає сплавам напів-Гайслера (НН) безліч надзвичайних властивостей. У цьому дослідженні ми вивчали структурні, електронні та термоелектричні властивості $ZrCoBi_{0.75}Z_{0.25}$ (Z = P, As, Sb), використовуючи теорію функціоналу щільності (DFT). Розрахунки проводились паралельно з використанням метода повного потенціалу лінеаризованої розширеної плоскої хвилі (FP-LAPW), який був реалізований в коді WIEN2k. Термоелектричні властивості були прогнозовані за допомогою напівкласичної теорії транспорту Больцмана, яка була реалізована в коді Больцтрапа. Отримані результати для зонної структури та щільностей станів підтверджують напівпровідникову (SC) природу трьох сполук із непрямим проміжком енергетичної зони, який становить близько 1 еВ. Основні термоелектричні параметри, такі як коефіцієнт Зеебека, теплопровідність, електропровідність та порівняльний показник якості, були оцінені для температур від нуля до 1200К. Позитивні значення коефіцієнта Зеебека (S) підтверджують, що $ZrCoBi_{0.75}Z_{0.25}$ (x = 0 та 0,25) є SC типу р. При температурі навколишнього середовища ZrCoBi_{0.75}P_{0.25} демонструє значну величину (S), яка складає 289 µV/K, що означає покращення на 22% порівняно з нелегованим ZrCoBi, а також демонструє зменшення теплопровідності на 54% (k/τ). Нелегований ZrCoBi має найнижче значення ZT при будь-яких температурах, і, при заміщенні атома вісмуту одним із sp-елементів (P, As, Sb), одночасне покращення k/τ та S призводить до досягнення максимальних значень порівняльного показника якості (ZT) \sim 0,84, отриманого при 1200 К для трьох легованих сполук.

КЛЮЧОВІ СЛОВА: першооснова, GGA, легований напівпровідник, коефіцієнт Зеебека, порівняльний показник якості